

Singly generated radical operator algebras

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ABSTRACT. We examine two nonselfadjoint operator algebras: the weighted shift algebra, and the Volterra operator algebra. In both cases, the operator algebra is the norm closure of the polynomials in the operator norm. In the case of the weighted shift algebra, the existence of a gauge action allows us to apply Fourier analysis to study the ideals of the algebra. In the case of the Volterra operator algebra, there is no gauge action, and other methods are needed to study the norm structure and the ideals.

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1. Introduction

Here we consider commutative operator algebras, which need not be self-adjoint. In the case of semi-simple commutative operator algebras, the Gelfand theory provides a complete description. At the other extreme, there is at this point no comprehensive theory of commutative radical operator algebras. This paper deals primarily with two types of commutative radical operator algebras: namely, those generated by weighted shifts, and the one generated by the Volterra integral operator.

Given a bounded linear operator T on a complex Hilbert space H , there are various topologies in which one take the closure of the polynomials in T to form an operator algebra. In this paper, we deal with the operator-norm closure. Thus, by the operator algebra \mathcal{A}_T we mean the operator-norm closure of the linear subspace of $\mathcal{B}(H)$ generated by $\{T, T^2, T^3, \dots\}$.

If T is a bounded linear operator on the Hilbert space H , then by definition the operator algebra \mathcal{A}_T is completely isometrically represented on H , and

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$C^*(T)$, the C^* -algebra generated by T in $\mathcal{B}(H)$, is a C^* -cover. The coordinate-free study of \mathcal{A}_T would include, say, the determination of the closed ideals of \mathcal{A}_T , rather than the invariant subspaces arising from the action of \mathcal{A}_T on the Hilbert space H , or, say, the existence of certain automorphisms, which is a property of the abstract operator algebra and not a particular representation. The coordinate-free study of operator algebras was stimulated by [1], which gave internal ‘matrix-norm’ conditions for a Banach algebra to be an operator algebra. Our approach here is necessarily a hybrid, as most properties of the operator algebra can only be deduced from the given representation, at least with the tools we have available.

One automorphism that has proved fruitful in the C^* -theory is the gauge automorphism. For example, this is useful in proving the simplicity of the Cuntz algebras \mathcal{O}_n (e.g., [3], Theorem V.4.6). However, gauge actions have been employed in nonselfadjoint operator algebras as well ([5]). For an operator $T \in \mathcal{B}(H)$, we say that the operator algebra \mathcal{A}_T admits a gauge action if the map $T \mapsto zT$ ($|z| = 1$) extends to an isometric isomorphism of \mathcal{A}_T . (See Definition 2.1.) The existence of a gauge action on \mathcal{A}_T allows for the application of Fourier analysis on the elements $S \in \mathcal{A}_T$, which in turn has application to the ideal structure of the algebra. For some operators $T \in \mathcal{B}(H)$ the associated operator algebra \mathcal{A}_T will admit a gauge action, while others will not. We show that if T is a weighted shift operator, then \mathcal{A}_T admits a gauge action. However if V is the Volterra integral operator, then \mathcal{A}_V fails to admit a gauge action. This distinction implies that \mathcal{A}_T and \mathcal{A}_V are not isomorphic as operator algebras. (Corollary 4.7)

Section 2 provides some background results regarding gauge actions and applications of gauge automorphisms to singly generated operator algebras \mathcal{A}_T , and some basic examples of operators T for which the associated algebra \mathcal{A}_T either does, or does not, admit a gauge action.

In Section 3 we consider operator algebras generated by weighted shift operators T . In addition to the operator norm on \mathcal{A}_T , there is a norm arising from a cyclic and separating unit vector for \mathcal{A}_T . But in general, the Hilbert space norm and the operator norm are inequivalent. (Remark 3.6) But if the weight sequence is square summable, then the two norms are equivalent. (Proposition 3.7) Proposition 3.5 gives a sufficient condition for an element S in the unit ball of \mathcal{A}_T to be an extreme point. In particular, the normalized powers of T , $T_n = \frac{1}{\|T^n\|} T^n$ are extreme points of the unit ball of \mathcal{A}_T . Under the same conditions on the weights, there is an isomorphism of the lattice of closed ideals of \mathcal{A}_T and closed invariant subspaces. (Proposition 3.10) We give two results describing which elements $S \in \mathcal{A}_T$ generate gauge-invariant ideals. We conclude this section showing that the operator algebra \mathcal{A}_T is a (nonunital) integral domain.

The final section of the paper deals with \mathcal{A}_V , the operator algebra generated by the classical Volterra integral operator on $L^2[0, 1]$. The closure of the polynomials in V in the strong operator topology turns out to be the commutant of

V , and hence corresponds also to the weak and weak* closed algebras generated by V . ([2] Theorem 5.10) Another weakly closed algebra associated with V is $\text{Alg}(\text{Lat } V)$, the weakly closed algebra of operators in $\mathcal{B}(L^2[0, 1])$ which leaves the lattice of subspaces $\text{Lat } V$ invariant. This algebra, which contains the commutant $\{V\}^c$, is non-commutative ([2] Theorem 5.12). The operator norm closed algebra generated by V , \mathcal{A}_V by contrast is composed of operators which share important properties of V : any $S \in \mathcal{A}_V$ is quasinilpotent and compact. Furthermore, given $S \in \mathcal{A}_V$ there is a measurable function f on $[0, 1]$, integrable over compact subsets of $[0, 1)$, such that if $\rho \in L^2[0, 1]$,

$$S\rho(x) = \int_0^x f(x-t)\rho(t) dt \text{ for almost all } x \in [0, 1]. (\text{Theorem 4.1})$$

However, f need not be integrable over $[0, 1]$, as shown in Example 4.2. If f is in $L^1[0, 1]$, then the L^1 norm of f dominates the operator norm $\|S\|$.

While Theorem 4.1 allows us to represent an arbitrary $S \in \mathcal{A}_V$ as a operator defined by a kernel, it does not provide another tool to calculate or estimate the norm. Even in the case of polynomials of low degree in V little is known. Remarkably, the recent paper [12] appears to be the first to have obtained an exact value for $\|V^2\|$, expressed as the solution to a transcendental equation ([12], Corollary 3.2¹). Their computations are limited to polynomials of degree 2 in V . As a consequence of this lack of computational tools, we are not able to make any assertions as to the extreme points of the unit ball of \mathcal{A}_V , as we did for the unit ball of the radical weighted shift algebra.

In [11] it is shown that the nilpotent elements are dense in \mathcal{A}_V . We obtain the same result here as a consequence of Theorem 4.1. The last result is an extension of Titchmarsh's theorem on zero divisors of $L^1[0, 1]$ to the operator algebra \mathcal{A}_V (Corollary 4.14).

2. Background, notation and examples

2.1. Gauge actions and Fourier analysis on singly generated algebras.

Let T be a bounded linear operator on a complex Hilbert space H , and \mathcal{A}_T the operator algebra in $\mathcal{B}(H)$ which is the operator norm closure of the polynomials in T which vanish at the origin. Let $\mathbb{T} = \{z \in \mathbb{C} : |z| = 1\}$.

Definition 2.1. Let $\text{Aut}(\mathcal{A}_T)$ denote the group of isometric automorphisms of \mathcal{A}_T . We say that \mathcal{A}_T admits a *gauge action* if there exists a continuous homomorphism $\gamma : \mathbb{T} \rightarrow \text{Aut}(\mathcal{A}_T)$ such that $\gamma_\lambda(T) = \lambda T$ ($\lambda \in \mathbb{T}$). By a *continuous homomorphism* we mean that for each $\lambda_0 \in \mathbb{T}$ and $S \in \mathcal{A}_T$,

$$\|\gamma_\lambda(S) - \gamma_{\lambda_0}(S)\| \rightarrow 0 \text{ as } \lambda \rightarrow \lambda_0 \text{ in } \mathbb{T}.$$

Remark 2.2. In the examples of gauge actions that arise here, the gauge automorphisms are completely isometric.

¹ $\|V^2\| = \eta_0^{-2}$, where η_0 is the least positive solution η to the equation $\cosh(\eta) \cos(\eta) = -1$. I recall many years ago hearing that Paul Halmos had obtained an expression for $\|V^2\|$ as the solution of a transcendental equation, but cannot find a reference for it.

Assume that \mathcal{A}_T admits a gauge action. Then if p is any polynomial with $p(0) = 0$, $\gamma_\lambda(p(T)) = p(\lambda T)$. Since, by definition, the algebra \mathcal{A}_T is the norm closure of such polynomials in T , it follows that the action of γ_λ on polynomials in T determines the action of γ_λ on \mathcal{A}_T .

Now if p is a polynomial, $p(z) = \sum_{j=1}^n a_j z^j$, then

$$\hat{p}(k)T^k = \int_{\mathbb{T}} \gamma_\lambda(p(T)) \lambda^{-k} d|\lambda| = \begin{cases} a_k T^k & \text{if } 1 \leq k \leq n \\ 0 & \text{otherwise} \end{cases}$$

Thus, for $S \in \mathcal{A}_T$,

$$\hat{S}(k)T^k = \int_{\mathbb{T}} \gamma_\lambda(S) \lambda^{-k} d|\lambda|$$

is well-defined. We say that $\hat{S}(k) \in \mathbb{C}$ is the k^{th} Fourier coefficient of S .

Lemma 2.3. *If \mathcal{A}_T admits a gauge action then $S \in \mathcal{A}_T$ is uniquely determined by its Fourier series.*

Proof. It is enough to prove that if $S \in \mathcal{A}_T$ is nonzero, then $\{\hat{S}(k)\}$ is not the zero sequence.

Suppose $S \neq 0$, and that $\hat{S}(k) = 0$ for all k . There is a continuous linear functional φ on \mathcal{A}_T for which $\varphi(S) \neq 0$, and hence the continuous function $f(\lambda) = \varphi(\gamma_\lambda(S))$ is nonzero. However,

$$\begin{aligned} \hat{f}(k) &= \int_{\mathbb{T}} \varphi(\gamma_\lambda(S)) \lambda^{-k} d|\lambda| \\ &= \varphi\left(\int_{\mathbb{T}} \gamma_\lambda(S) \lambda^{-k} d|\lambda|\right) \\ &= \varphi(\hat{S}(k)T^k) \\ &= \hat{S}(k)\varphi(T^k) \\ &= 0 \end{aligned}$$

This holds for $k = 1, 2, \dots$, but also for $k \leq 0$, since the Fourier coefficients $\hat{S}(k) = 0$ for all polynomials S and hence for all $S \in \mathcal{A}_T$.

This implies f is identically zero, which is a contradiction. \square

Just as with classical Fourier series, we associate with $S \in \mathcal{A}_T$ the formal series

$$S \sim \sum_{j=1}^{\infty} \hat{S}(j)T^j \quad (1)$$

We would like to construct a sequence of polynomials in T which converges to S in some sense. To this end, let φ be a continuous linear functional on \mathcal{A}_T and $p(z) = \sum_{j=1}^n a_j z^j$ a polynomial. Then

$$\widehat{\varphi(p(T))}(k) = \int_{\mathbb{T}} \varphi(\gamma_\lambda(p(T))) \lambda^{-k} d|\lambda| = a_k \varphi(T^k)$$

Since an arbitrary $S \in \mathcal{A}_T$ is a norm limit of polynomials, we have that

$$\widehat{\varphi(S)}(k) := \int_{\mathbb{T}} \varphi(\gamma_\lambda(S)) \lambda^{-k} d|\lambda| = a_k \varphi(T^k)$$

where $S \sim \sum_{j=1}^{\infty} a_j T^j$.

Proposition 2.4. *With notation as in the above paragraph, define the function $f : \mathbb{T} \rightarrow \mathbb{C}$, $f(\lambda) = \varphi(\gamma_\lambda(S))$. Then*

- 1 f is a continuous function on \mathbb{T} with $\hat{f}(n) = 0$ for $n \leq 0$.
- 2 The sequence of functions

$$\begin{aligned} s_n(\lambda) &= \sum_{j=1}^n \frac{n-j}{n} \hat{f}(j) \lambda^j \\ &= \sum_{j=1}^n \frac{n-j}{n} a_j \varphi(T^j) \lambda^j \quad (\lambda \in \mathbb{T}) \end{aligned}$$

converges uniformly in λ to f .

- 3 The sequence $\{S_n(\lambda) = \sum_{j=1}^n \frac{n-j}{n} a_j T^j \lambda^j\}$ converges weakly to $\gamma_\lambda(S)$, uniformly in $\lambda \in \mathbb{T}$.
- 4 There is a sequence R_n in the convex hull of the sequences $\{S_n(1) : n = 1, 2, \dots\}$ which converges in norm to S .

Proof. 1. By assumption, the map $\lambda \in \mathbb{T} \mapsto \gamma_\lambda(S) \in \mathcal{A}_T$ is norm continuous, and since φ is norm continuous, it follows that $f : \mathbb{T} \rightarrow \mathbb{C}$ is continuous. Now $\hat{S}(k) = 0$ for $k \leq 0$, so the same holds for f .

2. By Fejer's Theorem, the sequence of arithmetic means of the partial sums of the Fourier series for f converges uniformly to f on \mathbb{T} .

3. Since, for an arbitrary continuous linear functional φ , $\varphi(S_n(\lambda)) = s_n(\lambda)$, this is just a restatement of [2].

4. Follows from [3] by taking $\lambda = 1$ and applying the Hahn-Banach separation theorem. \square

Notation. The unitization of \mathcal{A}_T will be denoted $\tilde{\mathcal{A}}_T$.

At times it will be convenient to work in $\tilde{\mathcal{A}}_T$. The gauge action γ extends naturally to $\tilde{\mathcal{A}}_T$ with $\gamma_\lambda(I) = 1$. Of course for $S \in \tilde{\mathcal{A}}_T$, $\hat{S}(0)$ may be nonzero.

2.2. Nonselfadjoint operator algebras which admit a gauge action.

Example 2.5. Let \mathcal{M}_2 be the C^* -algebra of 2×2 matrices, with standard matrix units $e_{i,j}$, $1 \leq i, j \leq 2$. Let $T = e_{1,2}$. Then \mathcal{A}_T admits a gauge action. Indeed, since $T^2 = 0$, the operator space $\mathcal{A}_T = \mathbb{C} \cdot T$ is one-dimensional, and the map γ_λ is a linear map with $\gamma_\lambda(aT) = \lambda aT$, $a \in \mathbb{C}$.

To see that γ is completely isometric, it suffices to show that it extends to \mathcal{M}_2 . Define $U = e_{1,1} + \lambda e_{2,2}$. Then $\gamma_\lambda(A) = U^* A U$, $A \in \mathcal{M}_2$, $\lambda \in \mathbb{T}$ extends the action of γ on \mathcal{A}_T to the C^* -envelope, \mathcal{M}_2 .

Alternatively, we can invoke the description of \mathcal{M}_2 as the universal C^* -algebra generated by an operator T which is nilpotent of index 2 satisfying

$$T^*T + TT^* = I$$

$T \in \mathcal{B}(H)$ for some Hilbert space H , and I the identity in $\mathcal{B}(H)$. Since λT satisfies these same conditions for $\lambda \in \mathbb{T}$, it follows from the universal property that $T \mapsto \lambda T$ is automorphism of \mathcal{M}_2 .

Example 2.6. Let T be the multiplication operator on $L^2(\mathbb{T})$, $T\xi(z) = z\xi(z)$. The unital algebra $\tilde{\mathcal{A}}_T$ is the disc algebra $\mathcal{A}(\mathbb{D})$, and the algebra \mathcal{A}_T is the subalgebra of functions f satisfying $f \perp 1$, where 1 is the constant function in $L^2(\mathbb{T})$.

The gauge action γ_λ is given by

$$\gamma_\lambda(T)\xi(z) = \lambda z\xi(z).$$

Thus for $f \in \mathcal{A}(\mathbb{D})$, $\gamma_\lambda f(z) = f(\lambda z)$. This is isometric, even completely isometric. Indeed, the C^* -envelope of $\mathcal{A}(\mathbb{D})$ is $C(\mathbb{T})$, and the gauge action on the disc algebra is the restriction of the gauge action on $C(\mathbb{T})$, $\gamma_\lambda(f)(z) = f(\lambda z)$.

The Fourier series (as defined in equation 1) of $f \in \mathcal{A}(\mathbb{D})$ is the usual Fourier series of the function f .

Example 2.7. Let $\{S_1, \dots, S_d\}$ be isometries which satisfy the Cuntz relation

$$\sum_{j=1}^d S_j S_j^* = I.$$

Now if $i_1, \dots, i_n \in \{1, \dots, d\}$ and $\mu = (i_1, \dots, i_n)$ we write $S_\mu = S_{i_1} \dots S_{i_n}$ and $|\mu| = n$. Let \mathcal{A} be the Dirichlet algebra generated by the “monomials” $S_\mu S_\nu^*$ with $|\mu| \geq |\nu|$. Then \mathcal{A} is a nonself-adjoint subalgebra of the Cuntz algebra \mathcal{O}_d . Note that \mathcal{A} is invariant under the canonical gauge action on \mathcal{O}_d . Thus, \mathcal{A} admits a gauge action. The gauge action on this subalgebra of \mathfrak{D}_n was considered in [5].

Example 2.8. Let $\{S_1, \dots, S_d\}$ be the isometries of Example 2.7. If \mathcal{A} is the nonself-adjoint algebra generated by $\{S_1, \dots, S_d\} \subset \mathfrak{D}_n$, then \mathcal{A} admits a gauge action, since it is invariant under the canonical gauge action on \mathcal{O}_d . This algebra is known as Popescu’s noncommutative disc algebra.

Example 2.9. Let $\{S_1, \dots, S_d\}$ be as in Example 2.7. Here we assume that these operators are represented in some Hilbert space $\mathcal{B}(H)$. Choose one of the isometries, say S_1 , and let γ be the canonical gauge action on \mathcal{O}_d . Since $\gamma_\lambda(S_1) = \lambda S_1$, it follows that the subalgebra \mathcal{A}_{S_1} generated by S_1 of the Cuntz algebra \mathcal{O}_d is invariant under γ . Hence the gauge action on the Cuntz algebra \mathcal{O}_d restricts to a gauge action on \mathcal{A}_{S_1} .

Example 2.10. A variety of examples can be constructed as subalgebras of graph C^* -algebras which admit gauge actions. In this context one can obtain examples which are analogues of examples 2.7, 2.8 and 2.9, and where the generating isometries are replaced by Cuntz-Krieger partial isometries.

Let \mathcal{A} be an operator algebra, and $\mathfrak{A} = C_{env}^*(\mathcal{A})$ be its C^* -envelope. Then \mathcal{A}^* is an operator algebra defined as a subalgebra of \mathfrak{A} .

Proposition 2.11. *If γ is a gauge action on the operator algebra \mathcal{A} , then the adjoint algebra \mathcal{A}^* admits a gauge action, also denoted by γ defined by*

$$\gamma_\lambda(A^*) = (\gamma_\lambda(A))^* \quad \lambda \in \mathbb{T}, A \in \mathcal{A}$$

The proof is routine.

Proposition 2.12. *Every completely isometric automorphism of a unital operator algebra \mathcal{A} lifts to a $*$ -automorphism of the C^* -envelope $C_{env}^*(\mathcal{A})$, which fixes \mathcal{A} as a set.*

This is Proposition 10.1 of [4]. This tells us that if a unital operator algebra \mathcal{A} admits a gauge action γ , then each γ_λ extends to an automorphism, which we also denote by γ_λ , of the C^* -envelope, but does not immediately imply that the map $\lambda \in \mathbb{T} \mapsto \gamma_\lambda$ is continuous on the C^* -envelope.

2.3. Examples of operators in Hilbert space which do not admit a gauge action.

Example 2.13. Let $0 \neq P$ be a projection in $\mathcal{B}(H)$. As in Example 2.5 \mathcal{A}_P is one-dimensional, but in this case does not admit a gauge action. Indeed, since $P = P^2$, if γ were a gauge action on \mathcal{A}_P we would have

$$\lambda P = \gamma_\lambda(P) = \gamma_\lambda(P^2) = \gamma_\lambda(P)\lambda_\lambda(P) = \lambda^2 P, \quad \lambda \in \mathbb{T}$$

which is absurd.

Example 2.14. More generally, suppose that $T \in \mathcal{B}(H)$ is such that, for some $n > 1$, $0 \neq T^n$ and the set $\{T, T^2, \dots, T^n\}$ is linearly dependent. Then \mathcal{A}_T does not admit a gauge action.

Indeed, suppose to the contrary that \mathcal{A}_T admits a gauge action γ , and, choosing a dependence relation of minimal degree, we can assume that $a_1 T + \dots + a_m T^m = 0$, $m \leq n$ and $a_m \neq 0$.

Then

$$0 = \int_{\mathbb{T}} \gamma_\lambda \left(\sum_{k=1}^m (a_k T^k) \lambda^{-m} d|\lambda| \right) = a_m T^m$$

Since $T^m \neq 0$, it follows that $a_m = 0$, a contradiction.

Example 2.15. Let H be a Hilbert space with orthonormal basis $\{e_n\}_{n=1}^\infty$, and let $T \in \mathcal{B}(H)$ be the operator defined by $T e_n = \frac{1}{n} e_n$, $n \geq 1$. We claim that the operator algebra \mathcal{A}_T does not admit a gauge action.

Consider the operator $T - T^2 \in \mathcal{A}_T$. This is a compact, self-adjoint operator in $\mathcal{B}(H)$, so its norm is the maximum of the absolute values of the eigenvalues. $\|T - T^2\| = \|(T - T^2)e_2\|_2 = \frac{1}{4}$.

Suppose that \mathcal{A}_T admits a gauge action γ . Then $\gamma_\lambda(T - T^2) = \lambda T - \lambda^2 T^2$, so for $\lambda = -1$ we obtain $-T - T^2$. Computing $\| -T - T^2 \|$ we have $\| -T - T^2 \| = \|(-T - T^2)e_1\|_2 = 2$.

This is a contradiction, since by definition the gauge action is isometric on \mathcal{A}_T .

2.4. Gauge invariant Ideals in Operator algebras with gauge actions. Let \mathcal{A}_T be the operator algebra generated by an operator $T \in \mathcal{B}(H)$, and suppose \mathcal{A}_T admits a gauge action γ . A closed ideal $\mathcal{J} \subset \mathcal{A}_T$ is *gauge invariant* if, whenever $S \in \mathcal{J}$, then $\gamma_\lambda(S) \in \mathcal{J}$ ($\lambda \in \mathbb{T}$).

Proposition 2.16. *Let $\mathcal{J} \neq (0)$ be a gauge invariant ideal in \mathcal{A}_T . Then there exists $n \in \mathbb{N}$ such that $\mathcal{J} = \langle T^n \rangle$. That is, \mathcal{J} is the closed ideal in \mathcal{A}_T generated by T^n .*

Proof. Let $n = \inf\{k \geq 1 : \hat{S}(k) \neq 0 \text{ for some } S \in \mathcal{J}\}$. Thus, there exists $S \in \mathcal{J}$ with $\int_{\mathbb{T}} \gamma_\lambda(S) \lambda^{-n} d|\lambda| = a_n T^n \neq 0$. Since \mathcal{J} is closed and gauge invariant, $T^n \in \mathcal{J}$. It follows that any $S \in \mathcal{A}_T$ with Fourier series $S \sim \sum_{k=n}^{\infty} c_k T^k \in \mathcal{J}$. Thus, $\langle T^n \rangle \subset \mathcal{J}$. That is, the closed ideal generated by T^n is contained in \mathcal{J} .

On the other hand, let $S \in \mathcal{J}$. Then, by definition of n , S has Fourier series of the form $\sum_{k=n}^{\infty} c_k T^k$, so that $\mathcal{J} \subset \langle T^n \rangle$. \square

One ideal which is invariant under the gauge action is the Jacobson radical; indeed, it is invariant under all isometric automorphisms.

Corollary 2.17. *Let $T \in \mathcal{B}(H)$ be an operator such that \mathcal{A}_T admits a gauge action. Then either \mathcal{A}_T is semi-simple, or \mathcal{A}_T is radical.*

Proof. Let $\mathcal{J} \neq (0)$ denote the Jacobson radical of \mathcal{A}_T . Since the Jacobson radical is invariant under all isometric automorphisms, by Proposition 2.16 it follows that if the Jacobson radical is nonzero, there is an $n \in \mathbb{N}$ such that $\mathcal{J} = \langle T^n \rangle$. But if T^n is quasinilpotent, that is, has spectrum $\{0\}$, it follows from the Spectral Mapping Theorem that T has spectrum $\{0\}$. Hence, the ideal generated by T , which is \mathcal{A}_T is in the Jacobson radical. \square

Example 2.18. Here we note that it can happen that if $T \in \mathcal{B}(H)$ does not admit a gauge action, then we can have $(0) \neq \text{Rad}(\mathcal{A}_T) \neq \mathcal{A}_T$.

Let $H = H_1 \oplus H_2$ and $T = I_1 \oplus N$, where I_1 is the identity on H_1 and $N \in \mathcal{B}(H_2)$ is a nonzero nilpotent, with $N^2 = 0$. Let $p(z) = z - z^2$. Then $p(T) = 0 \oplus N \in \text{Rad}(\mathcal{A}_T)$, so that while \mathcal{A}_T is not a radical algebra, it has a non-trivial Jacobson radical.

The disc algebra $\mathcal{A}(\mathbb{D})$ has a rich lattice of ideals. ([6]) Not unexpectedly, there are few gauge invariant ideals.

Corollary 2.19. *If \mathcal{J} is a gauge invariant closed ideal of $\mathcal{A}(\mathbb{D})$, then (in the notation of Example 2.6) $\mathcal{J} = \langle z^n \rangle$ for some $n \in \mathbb{N}$.*

Proof. The conclusion follows immediately from Proposition 2.16. \square

3. Operator algebras generated by weighted shifts

In this section, T will denote a weighted shift operator. Let $\{e_n\}_{n \geq 0}$ be an orthonormal basis for the Hilbert space H , with $T e_n = a_n e_{n+1}$, $n \geq 0$, and

$a_n \neq 0$ for all n . Since T is bounded, the sequence $\{a_n\}$ is bounded, and $\|T\| = \sup_n |a_n|$.

We begin by showing that the operator algebra admits a gauge action.

Lemma 3.1. *Let T be as above. Then \mathcal{A}_T admits a gauge action.*

Proof. With $\{e_n\}_{n \geq 0}$ as above, define the unitary W_λ , $\lambda \in \mathbb{T}$, by $W_\lambda e_n = \lambda^n e_n$. Now

$$W_\lambda T W_\lambda^* e_n = W_\lambda T(\bar{\lambda}^n e_n) = \bar{\lambda}^n a_n W_\lambda e_{n+1} = \bar{\lambda}^n \lambda^{n+1} a_n e_{n+1} = \lambda T e_n$$

holds for any $n \geq 0$, and since the $\{e_n\}_{n \geq 0}$ form a basis, we have $W_\lambda T W_\lambda^* = \lambda T$. Thus the map $\lambda \in \mathbb{T} \mapsto W_\lambda T W_\lambda^* \in \mathcal{B}(H)$ is continuous, and so $\lambda \mapsto (W_\lambda T W_\lambda^*)^n = W_\lambda T^n W_\lambda^*$ is continuous, and hence $\lambda \mapsto W_\lambda p(T) W_\lambda^*$ for any polynomial p with $p(0) = 0$. Now if $S \in \mathcal{A}_T$ and $\epsilon > 0$ is given, there is a polynomial p with $\|p(T) - S\| < \epsilon/3$. Now let $\lambda_0 \in \mathbb{T}$ and $\delta > 0$ be such that if $|\lambda - \lambda_0| < \delta$, then $\|W_\lambda p(T) W_\lambda^* - W_{\lambda_0} p(T) W_{\lambda_0}^*\| < \epsilon/3$. Then

$$\begin{aligned} \|W_\lambda S W_\lambda^* - W_{\lambda_0} S W_{\lambda_0}^*\| &\leq \|W_\lambda (S - p(T)) W_\lambda^*\| + \\ &\quad \|W_\lambda p(T) W_\lambda^* - W_{\lambda_0} p(T) W_{\lambda_0}^*\| + \|W_{\lambda_0} (p(T) - S) W_{\lambda_0}^*\| \\ &< \epsilon/3 + \epsilon/3 + \epsilon/3 \end{aligned}$$

Thus $\gamma_\lambda(S) = W_\lambda S W_\lambda^*$ is a gauge action on \mathcal{A}_T . \square

Remark 3.2. We claim, furthermore, that the action is completely isometric. Now the C^* -algebra generated by T in $\mathcal{B}(H)$, $C^*(T)$, is a C^* -cover for \mathcal{A}_T , and the action of γ_λ on \mathcal{A}_T is the restriction to \mathcal{A}_T of the automorphism $S \in C^*(T) \mapsto \gamma_\lambda(S) := W_\lambda S W_\lambda^*$.

While it is clear that γ_λ is isometric on the C^* -cover, it is not obvious that the map $\lambda \in \mathbb{T} \mapsto \gamma_\lambda$ is continuous, since $\lambda \mapsto W_\lambda$ is not continuous.

It is more convenient to work with the unital algebras $\tilde{\mathcal{A}}_T, \tilde{\mathcal{A}}_T^*$. The C^* cover of $\tilde{\mathcal{A}}_T \subset \mathcal{B}(H)$ is the closure in $\mathcal{B}(H)$ of the union

$$\bigcup_{n=1}^{\infty} (\tilde{\mathcal{A}}_T^* \tilde{\mathcal{A}}_T)^n \subset \mathcal{B}(H)$$

Now since the action is continuous on $\tilde{\mathcal{A}}_T$ and $\tilde{\mathcal{A}}_T^*$, (Proposition 2.11) it is continuous on $(\tilde{\mathcal{A}}_T^* \tilde{\mathcal{A}}_T)^n$. And since it is isometric, it is thus continuous on the closure of the union.

Thus, the gauge action on \mathcal{A}_T is the restriction of a gauge action on a C^* -cover.

Lemma 3.3. *Let T be as in Lemma 3.1 If $S \in \mathcal{A}_T$, and*

$$S e_0 = \sum_{n=1}^{\infty} c_n e_n, \text{ then } c_n = \hat{S}(n) a_0 \dots a_{n-1}$$

Proof. Let W_λ ($\lambda \in \mathbb{T}$) be the family of unitary operators from Lemma 3.1, so that $W_\lambda e_n = \lambda^n e_n$. Let v be a linear combination of basis vectors. Since $\int_{\mathbb{T}} \lambda^{-n} W_\lambda v d|\lambda|$ is a multiple of e_n , it follows that,

$$\int_{\mathbb{T}} \lambda^{-n} W_\lambda v d|\lambda| = \langle v, e_n \rangle e_n$$

This holds for arbitrary vectors in H .

Suppose $Se_0 = \sum_{k=1}^{\infty} c_k e_k$. Then

$$\begin{aligned} c_n e_n &= \langle Se_0, e_n \rangle e_n \\ &= \int_{\mathbb{T}} W_\lambda (Se_0) \lambda^{-n} d|\lambda| \\ &= \int_{\mathbb{T}} W_\lambda S W_\lambda^* e_0 \lambda^{-n} d|\lambda| \\ &= \left(\int_{\mathbb{T}} W_\lambda S W_\lambda^* \lambda^{-n} d|\lambda| \right) e_0 \\ &= \left(\int_{\mathbb{T}} \gamma_\lambda(S) \lambda^{-n} d|\lambda| \right) e_0 \\ &= \hat{S}(n) T^n e_0 \end{aligned}$$

where we have used that $W_\lambda^* e_0 = e_0$.

Thus, $c_n e_n = \hat{S}(n) T^n e_0$, so that $c_n = \hat{S}(n)(a_0 a_1 \cdots a_{n-1})$. \square

Remark 3.4. Let T be as above, and $\tilde{\mathcal{A}}_T$ the unitization of \mathcal{A}_T . Then the vector e_0 is a cyclic and separating vector for $\tilde{\mathcal{A}}_T$. That it is cyclic is clear, for if v is any finite linear combination of basis vectors $v = \sum_{n=0}^N c_n e_n$, let p be the polynomial $p(z) = \sum_{n=0}^N \frac{c_n}{a_0 \cdots a_{n-1}} z^n$ (where the empty product is defined to be 1), then $p(T)e_0 = v$.

That e_0 is separating is also straightforward. First note that if $S \in \tilde{\mathcal{A}}_T$, there is a sequence of polynomials $\{p_n\}$ with $\{p_n(T)\}$ converging to S in the norm of $\tilde{\mathcal{A}}_T$, so that by definition of the norm, $p_n(T)v \rightarrow Sv$ for every $v \in H$ and in particular for $v = e_0$. By Proposition 2.4 these polynomials can be taken to be convex combinations of Fejer polynomials, so that $\hat{p}_n(k) \rightarrow \hat{S}(k)$ for every $k = 0, 1, \dots$. So if $S \neq 0$, there is some k with $\hat{S}(k) \neq 0$, and so $Se_0 = \sum_{j=0}^{\infty} \hat{S}(j) a_0 \cdots a_{j-1} e_j \neq 0$.

It is natural to ask for a description of the extreme points of the unit ball of \mathcal{A}_T . While that seems out of reach in our context, a sufficient condition is at hand.

Proposition 3.5. *Let the weights of T satisfy $|a_1| \geq |a_2| \geq \cdots$ and let $\mathcal{B} = \{S \in \mathcal{A}_T : \|S\| \leq 1\}$ be the closed unit ball. Then $S \in \mathcal{B}$ is an extreme point of \mathcal{B} if $\|S\| = \|Se_0\|_2 = 1$. In particular, the elements $T_n := \frac{1}{\|T^n\|} T^n$ are extreme points of \mathcal{B} .*

Proof. As noted in Remark 3.4, the map $S \in \mathcal{A}_T \mapsto \|Se_0\|_2$ is a norm on \mathcal{A}_T , satisfying $\|Se_0\|_2 \leq \|S\|$, and thus the map $S \mapsto Se_0$ maps the unit ball of \mathcal{A}_T into the unit ball of the Hilbert space H . Since every unit vector in Hilbert space is an extreme point of the unit ball in H , it follows that if $\|S\| = \|Se_0\|_2 = 1$, then Se_0 is extreme in the unit ball of H , and *a fortiori* S is extreme in the unit ball of \mathcal{A}_T . In particular, the condition on the weights $|a_1| \geq |a_2| \geq \dots$ guarantees that the monomials T^n assume their norm at e_0 , hence the normalized monomials T_n are extreme points of the unit ball of \mathcal{A}_T . \square

Remark 3.6. If the weights satisfy $|a_0| \geq |a_1| \geq |a_2| \geq \dots$ then, as noted in the proof of Proposition 3.5, the norms $\|T^n\|$, $\|T^n e_0\|_2$ coincide. However, it need not be the case that the two norms coincide, or even are equivalent, on the operator algebra \mathcal{A}_T .

To see this, let the weights satisfy $a_0 = a_1 = \dots = 1$, and take $p_n(T) = T + T^2 + \dots + T^n$ ($n \in \mathbb{N}$), and let $v_n = \frac{1}{\sqrt{n}}(e_0 + e_1 + \dots + e_{n-1})$. One calculates that

$$\|p_n(T)v_n\|_2 = \frac{\sqrt{2n^2 + 1}}{\sqrt{3}} \text{ while } \|p_n(T)e_0\|_2 = \sqrt{n}.$$

Thus,

$$\frac{\|p_n(T)\|}{\|p_n(T)e_0\|_2} \geq \frac{\sqrt{2}}{\sqrt{3}} \sqrt{n}$$

so the norms are inequivalent on \mathcal{A}_T .

While the operator norm is not in general equivalent to the norm $S \mapsto \|Se_0\|_2$ on \mathcal{A}_T , under certain restrictions the two norms are equivalent.

Proposition 3.7. *Let $\{a_n\}_{n \geq 0}$ be a sequence satisfying $|a_0| \geq |a_1| \geq |a_2| \geq \dots$ with $\sum_{n=0}^{\infty} |a_n|^2 := M^2 < \infty$, and $a_n \neq 0$ for all n . Then the operator norm on \mathcal{A}_T is equivalent to the norm $S \mapsto \|Se_0\|_2$.*

Proof. Let $p(z) = \sum_{j=1}^r c_j z^j$. Then

$$\begin{aligned} \|p(T)e_k\|_2^2 &= \sum_{j=1}^r |c_j|^2 |a_k a_{k+1} \dots a_{k+j-1}|^2 \\ &\leq \sum_{j=1}^r |c_j|^2 \left(\frac{|a_k|}{|a_0|}\right)^2 |a_0 a_1 \dots a_{j-1}|^2 \\ &\leq \left(\frac{|a_k|}{|a_0|}\right)^2 \|p(T)e_0\|_2^2 \end{aligned}$$

Now let v be a unit vector which is a finite linear combination of basis vectors, so $v = \sum_{\ell=0}^N \beta_\ell e_\ell$ with $\sum_{\ell=0}^N |\beta_\ell|^2 = 1$. Thus

$$\|p(T)v\|_2 \leq \sum_{\ell=0}^N |\beta_\ell| \|p(T)e_\ell\|_2$$

$$\begin{aligned}
&\leq \sum_{\ell=0}^N |\beta_\ell| \frac{|a_\ell|}{|a_0|} \|p(T)e_0\|_2 \\
&\leq \frac{1}{|a_0|} \|p(T)e_0\|_2 \left(\sum_{\ell=0}^N |\beta_\ell| |a_\ell| \right) \\
&\leq \frac{1}{|a_0|} \|p(T)e_0\|_2 \left(\sum_{\ell=0}^N |\beta_\ell|^2 \right)^{\frac{1}{2}} \left(\sum_{k=0}^N |a_k|^2 \right)^{\frac{1}{2}} \\
&\leq \frac{1}{|a_0|} \|p(T)e_0\|_2 \left(\sum_{\ell=0}^N |\beta_\ell|^2 \right)^{\frac{1}{2}} \left(\sum_{k=0}^{\infty} |a_k|^2 \right)^{\frac{1}{2}} \\
&\leq \frac{M}{|a_0|} \|p(T)e_0\|_2
\end{aligned}$$

Now since $\|p(T)v\|_2 \leq \frac{M}{|a_0|} \|p(T)e_0\|_2$ for a dense set of unit vectors v , it follows that $\|p(T)\| \leq \frac{M}{|a_0|} \|p(T)e_0\|_2$. Finally, since we can approximate an arbitrary $S \in \mathcal{A}_T$ by polynomials in T , so we conclude that $\|S\| \leq \frac{M}{|a_0|} \|Se_0\|_2$ for all $S \in \mathcal{A}_T$. \square

As a result of the equivalence of the two norms, several results follow immediately.

Corollary 3.8. *Let the weighted shift T be as in Proposition 3.7. Then for $S \in \mathcal{A}_T$, the partial sums of the Fourier series,*

$$\sum_{k=1}^n \hat{S}(k) T^k$$

converge in norm to S .

Corollary 3.9. *Let the weighted shift T be as in Proposition 3.7. Then the sequence*

$$p_n(T) := \sum_{k=1}^n c_k T^k$$

converges in norm to an element $S \in \mathcal{A}_T$ if and only if

$$\sum_{k=1}^{\infty} |c_k|^2 |a_0 \cdots a_{k-1}|^2 < \infty$$

Proposition 3.7 not only tells us that the operator norm is equivalent to a Hilbert space norm, but gives a mapping

$$\mathcal{F} : \mathcal{A}_T \rightarrow H, S \mapsto Se_0$$

which maps \mathcal{A}_T onto the closed subspace H_1 spanned by the basis vectors $e_n : n \geq 1$. One can also define $\tilde{\mathcal{F}} : \tilde{\mathcal{A}}_T \rightarrow H$ by $S \mapsto Se_0$. It is easy to see how to adapt Proposition 3.7 to the unital algebra $\tilde{\mathcal{A}}_T$. Note that the unital algebra

$\tilde{\mathcal{A}}_T$ maps onto H . Since \mathcal{F} is a Banach space isomorphism, it gives a one-to-one map of closed subspaces of \mathcal{A}_T to closed subspaces of H_1 , and similarly $\tilde{\mathcal{F}}$ maps closed subspaces of $\tilde{\mathcal{A}}_T$ onto closed subspaces of H . Furthermore

Proposition 3.10. *Suppose the weights of T satisfy the conditions of Proposition 3.7.*

- (1) *The map \mathcal{F} is an isomorphism of the lattice of closed ideals of \mathcal{A}_T onto the lattice of closed T -invariant subspaces of H_1 .*
- (2) *The map $\tilde{\mathcal{F}}$ is an isomorphism of the lattice of closed ideals of $\tilde{\mathcal{A}}_T$ onto the lattice of closed T -invariant subspaces of H .*

Proof. We prove only the second statement.

Let us first observe that the map $\tilde{\mathcal{F}}$ is T -equivariant. That is, $T\tilde{\mathcal{F}}(S) = \tilde{\mathcal{F}}(TS)$, $S \in \tilde{\mathcal{A}}_T$. Indeed, this follows immediately from the definition of $\tilde{\mathcal{F}}$. We claim that a closed subspace $\mathcal{J} \subset \tilde{\mathcal{A}}_T$ is a closed ideal if and only if it is T -invariant. Clearly, if \mathcal{J} is an ideal in $\tilde{\mathcal{A}}_T$, then it is T -invariant. On the other hand, if a closed subspace $\mathcal{J} \subset \tilde{\mathcal{A}}_T$ is T -invariant, then it is invariant under multiplication by any polynomial in T . Let $S \in \mathcal{J}$ and $R \in \tilde{\mathcal{A}}_T$. If $\{p_n\}$ is a sequence of polynomials such that $\{p_n(T)\}$ converges in norm to R , then $p_n(T)S$ converges in norm to RS . Thus, \mathcal{J} is a closed ideal.

Now clearly the map $\tilde{\mathcal{F}}$ maps closed subspaces of $\tilde{\mathcal{A}}_T$ to closed subspaces of H , and since $\tilde{\mathcal{F}}$ is T -equivariant, it is an isomorphism of the lattice of closed T -invariant subspaces of $\tilde{\mathcal{A}}_T$ onto the lattice of closed T -invariant subspaces of H . But as shown above, the closed T -invariant subspaces of $\tilde{\mathcal{A}}_T$ are exactly the closed ideals. \square

We know from Proposition 2.16 that the gauge invariant ideals in the algebra \mathcal{A}_T generated by a unilateral weighted shift T are all of the form $\langle T^k \rangle$. Given an element $S \in \mathcal{A}_T$, one can ask when the ideal $\langle S \rangle$ is of the form $\langle T^k \rangle$ for some $k \in \mathbb{N}$.

Corollary 3.11. *Let the weighted shift T be as in Proposition 3.7, and let $S \in \mathcal{A}_T$. Suppose $\hat{S}(j) = 0$, $j = 1, \dots, k-1$ and $\hat{S}(k) \neq 0$ for some $k > 1$. Then the closed ideal $\langle S \rangle = \langle T^k \rangle$ if and only if the closed subspace generated by the vectors $Se_0, TSe_0, T^2Se_0, \dots$ contains the basis vector e_k .*

Proof. The condition $\hat{S}(j) = 0$ for $j = 1, \dots, k-1$ implies that $\langle S \rangle \subset \langle T^k \rangle$. Indeed, there is a sequence of polynomials $\{p_n\} \subset \langle T^k \rangle$ converging in norm to S , hence $S \in \langle T^k \rangle$, and so the closed ideal $\langle S \rangle \subset \langle T^k \rangle$. Applying the map $\mathcal{F} : \mathcal{A}_T \rightarrow H$, it follows that Se_0 is contained in the closed invariant subspace generated by the vector e_k .

By Proposition 3.10, in order for the two ideals to coincide, the corresponding subspaces under the map \mathcal{F} must coincide. Thus it is necessary and sufficient that the closed subspace generated by the vectors $Se_0, TSe_0, T^2Se_0, \dots$ equal the closed subspace generated by e_k, Te_k, T^2e_k, \dots , which is the subspace generated by the basis vectors $e_k, e_{k+1}, e_{k+2}, \dots$. Thus, if the closed subspace generated by

the vectors $Se_0, TSe_0, T^2Se_0, \dots$ contains the vector e_k , by invariance it contains the vectors e_{k+1}, e_{k+2}, \dots , and hence the two closed subspaces coincide. \square

The weights $\{a_n\}$ satisfying the conditions of Proposition 3.7 satisfy $\lim_n a_n = 0$, so that the weighted shift T is quasinilpotent, and hence the algebra \mathcal{A}_T is radical. The consequences of the Lemma mentioned so far did not make direct use of the fact that the elements of the algebra are all quasinilpotent. The following Proposition gives a different sort of criterion as to when an element $S \in \mathcal{A}_T$ generates an ideal of the form $\langle T^k \rangle$. Here we do not need to assume the equivalence of the operator norm to the norm $S \mapsto \|Se_0\|_2$, rather we need only assume that the weighted shift T is quasinilpotent.

Recall ([13]) that a necessary and sufficient condition for a weighted shift operator to be quasinilpotent is that $\lim_n \sup_k |a_{k+1} \cdots a_{k+n}|^{\frac{1}{n}} = 0$.

Proposition 3.12. *Let $\{a_n\}_{n \geq 0}$ be a sequence of nonzero weights such that the unilateral weighted shift operator $Te_n = a_n e_{n+1}$ is quasinilpotent, and hence the algebra \mathcal{A}_T is radical. Let $S \in \mathcal{A}_T$ be a nonzero element such that $\hat{S}(j) = 0$, $j < k$ and $\hat{S}(k) \neq 0$. If in the unital algebra $\tilde{\mathcal{A}}_T$, S factors as $S = T^k Q$ for some $Q \in \tilde{\mathcal{A}}_T$, then $\langle S \rangle = \langle T^k \rangle$. In particular, if S is a polynomial $\hat{S}(k)T^k + \cdots + \hat{S}(n)T^n$, then $\langle S \rangle = \langle T^k \rangle$.*

Proof. First observe that $\langle S \rangle \subset \langle T^k \rangle$. By Proposition 2.4 there is a sequence of polynomials $\{p_n(T)\} \subset \langle T^k \rangle$ converging to $\langle S \rangle$. Thus S belongs to the closed ideal $\langle T^k \rangle$, and hence $\langle S \rangle \subset \langle T^k \rangle$.

Now we prove the reverse containment. Since multiplying S by a nonzero constant does not change the ideal $\langle S \rangle$, we may assume that $\hat{S}(k) = 1$, so that if $S = T^k Q$, then $\hat{Q}(0) = 1$. Writing $Q = I - R$, we have that $I - R$ is invertible in $\tilde{\mathcal{A}}_T$ with inverse $I + R + R^2 + \cdots$. Indeed, since $R \in \mathcal{A}_T$ which is radical, given any $r > 0$, $\|T^n\| \leq r^n$ for $n \geq N_r$. Thus, $S(I - R)^{-1} = S + SR + SR^2 + \cdots$. Note that while $S(I - R)^{-1}$ is a product in the unital algebra, the sum $S + SR + SR^2 + \cdots$ is computed in \mathcal{A}_T , and equals T^k . It follows that $T^k \in \langle S \rangle$, and hence $\langle T^k \rangle \subset \langle S \rangle$.

Finally observe that if S is a polynomial in T , then the factorization $S = T^k Q$ is realizable in $\tilde{\mathcal{A}}_T$. \square

In [14], Theorem 3, Shields characterizes the commutant of a weighted shift in terms of formal power series. In particular, that implies that the commutant is an integral domain. It follows that the smaller algebra \mathcal{A}_T is also an integral domain, though in our case a non-unital integral domain. The existence of the gauge action on \mathcal{A}_T allows us to deduce the same result.

Proposition 3.13. *Let T be a weighted shift with nonzero weight sequence as in Lemma 3.1. Then \mathcal{A}_T is a non-unital integral domain.*

In particular, if T is a quasinilpotent weighted shift, then the nonzero elements of \mathcal{A}_T are quasinilpotent and not nilpotent.

Proof. Let $R, S \in \mathcal{A}_T$, be nonzero elements. From Lemma 2.3 we know that the Fourier coefficients of R are not all zero, and similarly for S . By Proposition 2.4 there is a sequence of polynomials $\{p_n\}$ (resp., $\{q_n\}$) which are convex combinations of Fejer polynomials, so that $\{p_n(T)\}$ converges to R (resp., $\{q_n(T)\}$ converges to S). In particular, it follows from the Fejer property that $\widehat{p}_n(k) \neq 0$ implies $\widehat{R}(k) \neq 0$ (resp., $\widehat{q}_n(k) \neq 0$ implies $\widehat{S}(k) \neq 0$). Furthermore, for all k , $\widehat{p}_n(k) \rightarrow \widehat{R}(k)$ (resp., $\widehat{q}_n(k) \rightarrow \widehat{S}(k)$) as $n \rightarrow \infty$.

Let

$$j_0 = \min\{j : \widehat{R}(j) \neq 0\} \text{ and } k_0 = \min\{k : \widehat{S}(k) \neq 0\}.$$

Now $\{p_n(T)q_n(T)\}$ converges in norm to RS , and so $\{\widehat{p}_n\widehat{q}_n(\ell)\}$ converges to $\widehat{RS}(\ell)$ for all $\ell \in \mathbb{N}$. If $\ell_0 = \min\{\ell : \widehat{p}_n\widehat{q}_n(\ell) \neq 0 \text{ for } n \text{ sufficiently large}\}$ then $\ell_0 = j_0 + k_0$ and $\widehat{p}_n\widehat{q}_n(\ell_0) = \widehat{p}_n(j_0)\widehat{q}_n(k_0)$. Hence

$$\widehat{RS}(\ell_0) = \lim_n \widehat{p}_n(j_0)\widehat{q}_n(k_0) \neq 0$$

so that $RS \neq 0$.

For the second statement, if T is quasnilpotent, then the elements of \mathcal{A}_T , are quasnilpotent, and the nonzero elements are not nilpotent as \mathcal{A}_T is an integral domain. \square

4. The Volterra operator algebra

Let V be the Volterra operator on $L^2[0, 1]$, given by $V\xi(x) = \int_0^x \xi(t) dt$. Then we know ([8]) that for $n \geq 0$,

$$V^{n+1}\xi(x) = \frac{1}{n!} \int_0^x (x-t)^n \xi(t) dt \quad (2)$$

Let f be any $L^2[0, 1]$ function and let V_f denote the operator on $L^2[0, 1]$ given by

$$V_f\xi(x) = \int_0^x f(x-t)\xi(t) dt$$

Observe this is bounded, since

$$\begin{aligned} |V_f\xi(x)| &= \left| \int_0^x f(x-t)\xi(t) dt \right| \\ &\leq \int_0^x |f(x-t)| |\xi(t)| dt \\ &\leq \left[\int_0^x |f(x-t)|^2 dt \right]^{\frac{1}{2}} \left[\int_0^x |\xi(t)|^2 dt \right]^{\frac{1}{2}} \\ &\leq \|f\|_2 \|\xi\|_2 \end{aligned}$$

Now, since any $L^2[0, 1]$ function f is the L^2 limit of a sequence of polynomials, $\{p_n\}$, we have that

$$|(V_f - V_{p_n})\xi(x)| \leq \|f - p_n\|_2 \|\xi\|_2$$

so that

$$\|V_f - V_{p_n}\| \rightarrow 0 \quad (3)$$

Let \mathcal{A}_V denote the Volterra operator algebra, by which we mean the operator norm closure of the polynomials p in V with $p(0) = 0$. We have just shown that $V_f \in \mathcal{A}_V$ if $f \in L^2[0, 1]$. We would like to characterize arbitrary $T \in \mathcal{A}_V$.

Theorem 4.1. (1) *Let $f \in L^1[0, 1]$. Then for $\rho \in L^2[0, 1]$, the function $(V_f)\rho(x) := \int_0^x f(x-t)\rho(t) dt \in L^2[0, 1]$, and*

$$\|(V_f)\rho\|_2 \leq \|f\|_1 \|\rho\|_2 \text{ and hence } \|V_f\| \leq \|f\|_1$$

- (2) $f \in L^1[0, 1] \mapsto \|V_f\|$ is a norm on $L^1[0, 1]$.
 (3) Let $T \in \mathcal{A}_V$ and $\{f_n\}$ a sequence of functions in $L^1[0, 1]$ such that $V_{f_n} \rightarrow T$ in \mathcal{A}_V . If $0 < x_0 < 1$, then the sequence $\{\int_0^{x_0} |f_n(x)| dx\}$ is bounded.
 (4) Let $0 < x_0 < 1$ and let $S_{x_0} = \{f : f \in L^1[0, 1], f(x) = 0 \text{ for } x_0 < x \leq 1\}$. Then on S_{x_0} the operator norm $f \mapsto \|V_f\|$ and the L^1 -norm $f \mapsto \int_0^1 |f|$ are equivalent.
 (5) Let $T \in \mathcal{A}_V$. Then there is a measurable function f on $[0, 1]$, integrable over compact subsets of $[0, 1]$, such that $T = V_f$.
 (6) If f, g are measurable functions on $[0, 1]$ such that $V_f, V_g \in \mathcal{A}_V$ and $\alpha \in \mathbb{C}$, then $V_{\alpha f} = \alpha V_f$ and $V_{f+g} = V_f + V_g$.

Proof. Let $\rho \in L^2[0, 1]$ and $f \in L^1[0, 1]$. For $x \in [0, 1]$ define the function f_x by

$$f_x(t) = \begin{cases} f(x-t) & \text{if } 0 \leq t \leq x \\ 0 & \text{if } x < t \leq 1 \end{cases}$$

$$\begin{aligned} |(V_f)\rho(x)| &= \left| \int_0^x f(x-t)\rho(t) dt \right| \\ &\leq \int_0^x |f(x-t)| |\rho(t)| dt \\ &\leq \int_0^1 (\sqrt{|f_x(t)|}) (\sqrt{|f_x(t)| |\rho(t)|}) dt \\ &\leq \left(\int_0^1 |f_x(t)| dt \right)^{\frac{1}{2}} \left(\int_0^1 |f_x(t)| |\rho(t)|^2 dt \right)^{\frac{1}{2}} \\ &\leq \|f_x\|_1^{\frac{1}{2}} \left(\int_0^1 |f_x(t)| |\rho(t)|^2 dt \right)^{\frac{1}{2}} \\ &\leq \|f\|_1^{\frac{1}{2}} \left(\int_0^1 |f_x(t)| |\rho(t)|^2 dt \right)^{\frac{1}{2}} \end{aligned}$$

Thus

$$\begin{aligned} \int_0^1 |(V_f)\rho(x)|^2 dx &\leq \|f\|_1 \int_0^1 \int_0^1 |f_x(t)| |\rho(t)|^2 dt dx \\ &\leq \|f\|_1 \int_0^1 \left(\int_0^1 |f_x(t)| dx \right) |\rho(t)|^2 dt \\ &\leq \|f\|_1 \|f\|_1 \|\rho\|_2^2 \\ &\leq (\|f\|_1 \|\rho\|_2)^2 \end{aligned}$$

Thus, $(V_f)\rho \in L^2[0, 1]$ and $\|(V_f)\rho\|_2 \leq \|f\|_1 \|\rho\|_2$. Thus $\|V_f\| \leq \|f\|_1$. To prove (2), note that by definition,

$$\|V_f\| = \sup\{ | \langle (V_f)\rho, \xi \rangle | : \|\rho\|_2 \leq 1, \|\xi\|_2 \leq 1 \}.$$

For $f \in L^1[0, 1]$, define

$$\rho(t) = \overline{\operatorname{sgn}(f(1-t))} := \begin{cases} \frac{\overline{f(1-t)}}{|f(1-t)|} & \text{if } f(1-t) \neq 0 \\ 0 & \text{otherwise} \end{cases}$$

and let $F(x) = \int_0^x f(x-t)\rho(t) dt$. Then F is continuous, and $F(1) = \|f\|_1$. Now let $\xi(x) = \operatorname{sgn}(F(x))$. Since $\|\rho\|_2, \|\xi\|_2 \leq 1$, we have

$$\|V_f\| \geq | \langle (V_f)\rho, \xi \rangle | = \int_0^1 |F(x)| dx > 0.$$

For (3), if the sequence of integrals is not bounded, then for all $x_0 \leq x \leq 1$, the values of $T(1)(x)$ are either $\pm\infty$ or undefined. But that contradicts that $T(1) \in L^2[0, 1]$.

For (4), if the space S_{x_0} is complete in the operator norm $f \mapsto \|V_f\|$, then since by part (1) $\|V_f\| \leq \|f\|_1$, it follows from a Corollary of the Open Mapping Theorem that the two norms are equivalent. Suppose that S_{x_0} is not complete in the operator norm. Then there is an element $T \in \mathcal{A}_V$, $\|T\| = 1$, and a sequence $\{f_n\} \subset S_{x_0}$ so that $\{V_{f_n}\}$ converges to T with $\int_0^{x_0} |f_n|$ unbounded. But that contradicts (3).

For (5), let $0 < x_0 < 1$ and define the projection P_{x_0} on $L^2[0, 1]$ by

$$P_{x_0}\rho(t) = \begin{cases} \rho(t) & \text{if } 0 \leq t \leq x_0 \\ 0 & \text{otherwise} \end{cases}$$

For $f \in L^1[0, 1]$, let $f_{x_0} \in S_{x_0}$ be defined by $f_{x_0}(t) = \begin{cases} f(t) & \text{if } 0 \leq t \leq x_0 \\ 0 & \text{otherwise} \end{cases}$

Observe that

$$P_{x_0}V_fP_{x_0} = P_{x_0}V_{f_{x_0}}P_{x_0}$$

Also note that $f \in S_{x_0} \mapsto \|P_{x_0} V_f P_{x_0}\|$ is a norm on S_{x_0} , weaker than the operator norm. That it is a norm follows from (2), with the interval $[0, 1]$ replaced by $[0, x_0]$.

Let $T \in \mathcal{A}_V$ and $\{f_n\}$ a sequence of functions in $L^1[0, 1]$ such that $\{V_{f_n}\}$ converges to T . Then

$$\{P_{x_0} V_{f_n} P_{x_0} = P_{x_0} V_{f_{n,x_0}} P_{x_0}\} \text{ converges to } P_{x_0} T P_{x_0}$$

In other words, the sequence $\{V_{f_{n,x_0}}\}$ converges in the norm defined in the previous paragraph. Since by (3) the integrals $\int_0^{x_0} |f_n|$ are bounded, the limit of the sequence $\{f_{n,x_0}\}$ with respect to this norm belongs to S_{x_0} . Thus, the limit of $\{V_{f_{n,x_0}}\}$ has the form V_g for some $g \in S_{x_0}$.

Now if $x_0 < y < 1$, then $\{f_{n,y}\}$ converges with respect to the norm $q \in S_y \mapsto \|P_y V_q P_y\|$ to a function h . In other words, $\{V_{f_{n,y}}\}$ converges in this norm to an operator V_h for some $h \in S_y$, and furthermore, the restriction of h to the interval $[0, x_0]$ coincides with g . Thus, if $x_0 < x_1 < x_2 < \dots < 1$ with $\lim_n x_n = 1$, then we obtain a sequence of functions f_{x_n} such that $f_{x_{n+1}}$ restricted to $[0, x_n]$ equals f_{x_n} on that interval. Thus we obtain a function f , measurable on $[0, 1]$, whose restriction to $[0, x_n]$ is equal to the restriction of f_{x_n} to $[0, x_n]$. Finally, the fact that f is integrable over compact subsets of $[0, 1)$ follows from the fact that all of the f_{x_n} are integrable.

To verify (6), let $\{p_n\}, \{q_n\}$ be sequences of polynomials such that $\{V_{p_n}\}, \{V_{q_n}\}$ converge to V_f, V_g respectively. Since $V_{p_n+q_n} = V_{p_n} + V_{q_n}$, and $V_{\alpha p_n} = \alpha V_{p_n}$, the conclusion follows. \square

In [9] G. Little and J. B. Reade prove an asymptotic estimate for the norm of powers of the Volterra operator V :

$$\lim_n n! \|V^n\| = \frac{1}{2}.$$

This result can be interpreted as an asymptotic estimate of the ratio of the operator norm to the L^1 -norm on the set of functions $f_n(x) = x^n$. Indeed, since $n! V^{n+1} = V_{f_n}$, we have

$$\frac{\|V_{f_n}\|}{\|f_n\|_1} = \frac{n! \|V^{n+1}\|}{\frac{1}{n+1}} = (n+1)! \|V^{n+1}\| \rightarrow \frac{1}{2}$$

as $n \rightarrow \infty$. Since the set of functions $\{f_n : n = 0, 1, \dots\}$ spans a dense subspace of the operator algebra \mathcal{A}_V , this result seems to suggest that the two norms may be equivalent.

However, it turns out that the two norms are not equivalent, as the following example shows.

Example 4.2. Let \mathcal{S} be the space of Lebesgue measurable functions f on $[0, 1]$ such that, for every $0 < x < 1$, $f|_{[0,x]} \in L^2[0, x]$. For $f \in \mathcal{S}$, define

$$\rho_x(t) = \begin{cases} \frac{1}{c(x)} \overline{f(x-t)}, & \text{if } c(x) \neq 0 \text{ and } t \leq x \\ 0, & \text{otherwise} \end{cases}$$

where $c(x) = [\int_0^x |f(x-t)|^2 dt]^{\frac{1}{2}}$. Then, clearly, for any function $\rho \in L^2[0, 1]$ of unit norm, $|\int_0^x f(x-t)\rho(t) dt| \leq \int_0^x f(x-t)\rho_x(t) dt$.

Thus, if $G(x) = (\int_0^x f(x-t)\rho_x(t) dt)^2$, and if f is not zero a.e. in the interval $[0, x]$, we have

$$\begin{aligned} G(x) &= \frac{1}{c(x)^2} \left(\int_0^x |f(x-t)|^2 dt \right)^2 \\ &= \int_0^x |f(x-t)|^2 dt \end{aligned}$$

Hence $\|V_f\| \leq [\int_0^1 G(x) dx]^{\frac{1}{2}}$.

Let $\mathcal{S}_1 = \{f \in \mathcal{S} : \int_0^1 \int_0^x |f|^2 < \infty\}$, and for

$$f \in \mathcal{S}_1, \text{ set } \|f\|_{\#} = \left[\int_0^1 \int_0^x |f|^2 \right]^{\frac{1}{2}}.$$

If the kernel k_f is defined by

$$k_f(x, t) = \begin{cases} f(x-t) & \text{if } t \leq x \\ 0 & \text{if } t > x \end{cases}$$

then the condition $f \in \mathcal{S}_1$ is equivalent to $k_f \in L^2([0, 1] \times [0, 1])$ and in that case $\|f\|_{\#} = \|k_f\|_2$. Thus $\|f\|_{\#}$ is the Hilbert-Schmidt norm of the operator V_f .

To show that \mathcal{A}_V properly contains $L^1[0, 1]$ it suffices to exhibit a function f with $\|f\|_{\#} < \infty$ but $f \notin L^1[0, 1]$.

Now let $f = \sum_{n=1}^{\infty} \frac{2^n}{n} \chi_{[1-2^{-(n-1)}, 1-2^{-n}]}$. Then $\int_0^1 f = \sum_{n=1}^{\infty} \frac{1}{n}$ diverges, so $f \notin L^1[0, 1]$. However, $\|f\|_{\#}^2 = \int_0^1 G$ is finite. To see this, view the area under the graph of G as divided into horizontal strips. The portion of the area between $G(1-2^{-(n-1)})$ and $G(1-2^{-n})$ is $\frac{3}{2} \cdot \frac{1}{n^2}$. Thus,

$$\int_0^1 G = \sum_{n=1}^{\infty} \frac{3}{2} \cdot \frac{1}{n^2} = \pi^2/4.$$

It follows that the operator norm $\|V_f\|$ is finite, and in fact $\|V_f\| \leq \pi/2$. If $f_n := f \chi_{[0, 1-2^{-n}]}$, then $f_n \in L^1[0, 1]$ and a calculation similar to the above

shows that $\|V_f - V_{f_n}\| \rightarrow 0$ as $n \rightarrow \infty$. Since $V_{f_n} \in \mathcal{A}_V$ and \mathcal{A}_V is by definition complete, $V_f \in \mathcal{A}_V$.

Thus, $V_f \in \mathcal{A}_V$, but $f \notin L^1[0, 1]$.

Remark 4.3. There is no simple relationship between the Hilbert Schmidt norm of V_f and $\|f\|_1$. We have just seen that the Hilbert-Schmidt norm of V_f can be finite and f not in $L^1[0, 1]$. On the other hand, if $f \in L^1[0, 1]$ is such that, for some $0 < x < 1$, $f|_{[0,x]} \notin L^2[0, x]$, then the L^1 norm of f is finite but the Hilbert Schmidt norm of V_f is not.

Example 4.4. If f is measurable on $[0, 1]$, the condition $k_f \in L^1([0, 1]^2)$ does not imply $V_f \in \mathcal{A}_V$. Let $f(x) = \frac{1}{(1-x)^{\frac{3}{2}}}$. Then

$$\|k_f\|_1 = \int_0^1 \left(\int_0^x f(t) dt \right) dx < \infty$$

To show V_f is unbounded, it is enough to find $\rho \in L^2[0, 1]$ such that $(V_f)\rho \notin L^2[0, 1]$. Take ρ to be the constant 1. Then

$$\|(V_f)\rho\|_2^2 = \int_0^1 \left| \int_0^x f(t) dt \right|^2 dx \text{ diverges.}$$

This shows that the conditions of Theorem 4.1 part (5) on a measurable function f to satisfy $V_f \in \mathcal{A}_V$ are necessary but not sufficient.

Remark 4.5. Let f be measurable on $[0, 1]$ and integrable over compact subsets of $[0, 1)$. One would like to have an ‘easily verifiable’ necessary and sufficient condition for V_f to belong to \mathcal{A}_V . One might conjecture that such a condition depends on the rate of growth of $\{\int_0^x |f|\}$ as $x \uparrow 1$. On the other hand, could it happen, for some f , that $V_f \in \mathcal{A}_V$, but $V_{|f|} \notin \mathcal{A}_V$?

4.1. Operator algebraic properties of \mathcal{A}_V . We turn now from a discussion of the norm of operators in \mathcal{A}_V to its properties as an algebra. Since the Volterra operator is quasinilpotent, the algebra \mathcal{A}_V is a commutative radical operator algebra. From the example of the weighted shift, we know that radical operator algebras \mathcal{A}_T can admit a gauge action. Does the same hold for \mathcal{A}_V ?

Proposition 4.6. \mathcal{A}_V does not admit a gauge action.

Proof. We will make use of the formula for V^{n+1} (cf equation 2).

Now by the Muntz-Szasz Theorem, the function $g(x) = x$ can be uniformly approximated in $[0, 1]$ by polynomials in $\{x^2, x^3, \dots\}$. Thus given $\epsilon > 0$, we can find a polynomial $p(x) = a_2x^2 + a_3x^3 + \dots + a_nx^n$ satisfying $\|g - p\|_\infty = \sup_{0 \leq x \leq 1} |x - p(x)| < \epsilon$.

Let $\rho \in L^2[0, 1]$ with $\|\rho\|_2 = 1$. Then

$$\|V^2(\rho) - \sum_{j=2}^n j! a_j V^{j+1}(\rho)\|_2 = \left\| \int_0^x [(x-t) - p(x-t)]\rho(t) dt \right\|_2$$

$$\begin{aligned}
&\leq \left\| \int_0^x |(x-t) - p(x-t)| |\rho(t)| dt \right\|_2 \\
&\leq \left\| \int_0^1 \epsilon |\rho(t)| dt \right\|_2 \\
&\leq \epsilon
\end{aligned}$$

Since this holds for all $\rho \in L^2[0, 1]$ of norm 1, it follows that

$$\left\| V^2 - \sum_{j=2}^n j! a_j V^{j+1} \right\| \leq \epsilon.$$

Now we assume that the operator algebra \mathcal{A}_V admits a gauge action, γ . Since by definition the Fourier coefficient $\widehat{V^2}(2) = 1$, and $(\sum_{j=2}^n j! a_j V^{j+1})(2) = 0$, and since the operation $S \in \mathcal{A}_V \mapsto \widehat{S}(2) = \int_{\mathbb{T}} \gamma_\lambda(S) \lambda^{-2} d|\lambda|$ is norm-decreasing, it follows that

$$\|V^2\| = \left\| \int_{\mathbb{T}} (\gamma_\lambda(V^2 - \sum_{j=2}^n j! a_j V^{j+1}) \lambda^{-2} d|\lambda| \right\| \leq \epsilon$$

which is absurd, since $\epsilon > 0$ was arbitrary.

Thus \mathcal{A}_V does not admit a gauge action. \square

If T is a quasinilpotent weighted shift, then the fact that the lattice $\text{Lat } T$ is discrete and $\text{Lat } V$ is continuous tells us that the two operators are not unitarily equivalent. But to show that the operator algebras \mathcal{A}_T , \mathcal{A}_V are not isomorphic requires another argument.

Corollary 4.7. *Let T be a quasinilpotent weighted shift. Then the radical operator algebras \mathcal{A}_T , \mathcal{A}_V are not completely isometrically isomorphic.*

Proof. Since \mathcal{A}_T admits a gauge action, and \mathcal{A}_V does not, they cannot be completely isometrically isomorphic. \square

Define a ‘convolution’ on elements of \mathcal{A}_V as follows: if f, g are measurable functions on $[0, 1]$ such that $V_f, V_g \in \mathcal{A}_V$, set

$$f * g(x) = \int_0^x f(x-t)g(t) dt \tag{4}$$

First observe that since the restrictions of f, g to the interval $[0, x]$ are integrable if $x < 1$, it follows that $f * g$ is well defined. Furthermore, a calculation shows that, for $\rho \in L^2[0, 1]$,

$$V_{f*g}(\rho) = V_f(V_g(\rho)) \text{ and hence } V_{f*g} = V_f V_g. \tag{5}$$

It is well known that the closed invariant subspaces of the Volterra operator V have the form $\{\xi \in L^2[0, 1] : \xi(t) = 0 \text{ for } 0 \leq t \leq x_0\}$, for $0 < x_0 < 1$. ([2], Theorem 5.5. Note that their notation differs from our: their V is V^* here.) The

same holds for the operator algebra, \mathcal{A}_V . This yields some information about the closed ideals of \mathcal{A}_V .

Proposition 4.8. *Let $0 < x_0 < 1$. The subspace $\mathcal{J}_{x_0} = \{V_f : f(t) = 0, \text{ for } 0 \leq t \leq x_0\}$ is a closed ideal of \mathcal{A}_V .*

Proof. Let $V_g \in \mathcal{J}_{x_0}$ and $V_f \in \mathcal{A}_V$. Since $V_f V_g = V_{f * g}$, and since we have $f * g(x) = \int_0^x f(x-t)g(t)dt$, it is clear that if $x \leq x_0$ then $f * g(x) = 0$. Thus, \mathcal{J}_{x_0} is an ideal.

The proof that \mathcal{J}_{x_0} is a closed ideal is similar to the proof of statement (2) of Theorem 4.1. Suppose $V_f \in \mathcal{A}_V$, $V_f \notin \mathcal{J}_{x_0}$. This implies that the restriction of f to the interval $[0, x_0]$ is not zero.

Let $F(x) = \int_0^x f(x-t)\rho(t)dt$, where

$$\rho(t) = \begin{cases} \text{sgn}(f(x_0 - t)), & \text{if } t \leq x_0 \\ 0, & \text{if } t > x_0 \end{cases}$$

Then F is continuous, and $F(x_0) = \int_0^{x_0} |f(x_0 - t)| dt > 0$. Let

$$\xi(t) = \begin{cases} \text{sgn}(F(x)), & \text{if } x \leq x_0 \\ 0, & \text{if } x > x_0 \end{cases}$$

Now let $V_g \in \mathcal{J}_{x_0}$ be arbitrary. Then

$$\langle (V_f - V_g)(\rho), \xi \rangle = \int_0^{x_0} |F(x)| dx \text{ is a positive constant.}$$

It follows that no sequence $\{V_{g_n}\} \subset \mathcal{J}_{x_0}$ converges to V_f . Hence \mathcal{J}_{x_0} is a closed ideal. \square

Remark 4.9. In Proposition 3.10 we established a one-to-one correspondence between closed ideals of the operator algebra \mathcal{A}_T of the weighted shift T , and invariant subspaces. If such a relationship were to exist for the Volterra operator algebra, then we would have a complete description of the closed ideals of \mathcal{A}_V .

Note that if an operator algebra \mathcal{A} is completely isometrically represented on a Hilbert space H with cyclic vector ξ_0 , then given an invariant subspace H_1 there is a closed ideal \mathcal{J} by $\mathcal{J} = \{a \in \mathcal{A} : a\xi_0 \in H_1\}$. In the other direction, there is no assurance that if \mathcal{J} is a closed ideal of \mathcal{A} , that the subspace $\mathcal{J} \cdot \xi_0$ is closed in H .

Lemma 4.10. *Let $\alpha, \beta \in (0, 1)$ and suppose f, g are measurable functions on $[0, 1]$ such that $V_f, V_g \in \mathcal{A}_V$, and*

- (1) *If f is supported on $[\alpha, 1]$ and g is supported on $[\beta, 1]$, then $f * g$ is supported on $[(\alpha + \beta), 1]$ if $\alpha + \beta < 1$, and $f * g = 0$ if $\alpha + \beta \geq 1$.*
- (2) *In particular, if $\beta = 1 - \alpha$, then $f * g = 0$.*

Proof. The proof is a straightforward computation, making use of the convolution formula, equation 4. \square

Corollary 4.11. *If f is a measurable function on $[0, 1]$ such that $V_f \in \mathcal{A}_V$ and f is supported on $[\alpha, 1]$, then $(V_f)^n = 0$ if $n\alpha \geq 1$.*

Proof. This follows from repeated application of Lemma 4.10 □

The following result appears as Proposition 2.13 of [11]. It is proved there using methods entirely different from those in this paper.

Corollary 4.12. *The nilpotent elements in the Volterra operator algebra \mathcal{A}_V are dense.*

Proof. Let $V_f \in \mathcal{A}_V$. Since the elements of the form V_g , $g \in L^1[0, 1]$ are dense in \mathcal{A}_V , given $\epsilon > 0$, let $g \in L^1[0, 1]$ satisfy $\|V_f - V_g\| < \epsilon$. There is $\delta > 0$ such that if $E \subset [0, 1]$ is measurable with $\mu(E) \leq \delta$, then $\int_E |g| \leq \epsilon$.

Define $h \in L^1[0, 1]$ by

$$h(x) = \begin{cases} 0, & \text{if } 0 \leq x < \delta \\ g(x), & \text{if } \delta \leq x \leq 1 \end{cases}$$

Then

$$\|V_f - V_h\| \leq \|V_f - V_g\| + \|V_g - V_h\| < \epsilon + \|g - h\|_1 < 2\epsilon$$

Observe that the second inequality above makes use of part (1) of Theorem 4.1, since

$$\|V_g - V_h\| = \|V_{g-h}\| \leq \|g - h\|_1$$

Let $n \in \mathbb{N}$ satisfy $n\delta \geq 1$. It follows from Corollary 4.11 that V_h is nilpotent in \mathcal{A}_V . □

Remark 4.13. $L^1[0, 1]$ is a Banach algebra under the convolution $f * g(x) = \int_0^x f(x-t)g(t) dt$. Hence this Banach algebra is dense in \mathcal{A}_V , and the convolution is the restriction of that in \mathcal{A}_V to L^1 .

One version of the Titchmarsh convolution is: if $f, g \in L^1[0, 1]$ and $f * g = 0$ then there is an $\alpha \in [0, 1]$ such that $\text{supp } f \subset [\alpha, 1]$ and $\text{supp } g \subset [1 - \alpha, 1]$. ([2] Problem 5.4 and [8] Theorem 10, Sec. 38.3)

The next result shows that the Titchmarsh convolution theorem also holds in the larger algebra \mathcal{A}_V . In particular, for real-valued functions the converse of statement (2) of Lemma 4.10 holds.

Corollary 4.14. *Let f, g be real-valued measurable functions on $[0, 1]$ such that $V_f, V_g \in \mathcal{A}_V$. Then $V_{f*g} = 0$ if and only if there exists $\alpha \in [0, 1]$ such that $\text{supp } f \subset [\alpha, 1]$ and $\text{supp } g \subset [1 - \alpha, 1]$.*

Proof. One direction follows from Lemma 4.10.

If $V_{f*g} = V_f V_g = 0$, then $V(V_f V_g) = V_F V_g = 0$ where $F(x) = \int_0^x f(t) dt$. It follows that $V_F V_g \xi = 0$ for all real-valued $\xi \in L^2[0, 1]$. Thus $\int_0^x F(x-t)(V_g \xi)(t) dt = 0$ for a.a. $x \in [0, 1]$.

As $F = V_f(1)$ and $V_g \xi$ are both in $L^2[0, 1] \subset L^1[0, 1]$, it follows from the classical Titchmarsh theorem that there exists $\alpha \in [0, 1]$ such that $F = 0$ in $[0, \alpha]$ and $\int_0^x g(x-t)\xi(t) dt = 0$ for $x \in [0, 1 - \alpha]$. Thus $f = 0$ in $[0, \alpha]$.

Suppose $g \neq 0$ in $[0, 1 - \alpha]$. Thus for some $x_0 \in [0, 1 - \alpha]$, $\int_0^{x_0} |g| dt \neq 0$. Define $\xi(t) = \text{sgn}(g(x_0 - t))$. Then $\int_0^x g(x - t)\xi(t) dt$ is a nonzero continuous function in the interval $[0, 1 - \alpha]$, a contradiction. \square

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